

5. The Gradient Operator

A brief review is provided here for the gradient operator $\bar{\nabla}$ in both Cartesian and orthogonal non-Cartesian coordinate systems.

Sections in this Chapter:

5.01 Gradient, Divergence, Curl and Laplacian (Cartesian)

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5.01 Gradient, Divergence, Curl and Laplacian (Cartesian)

Let z be a function of two independent variables (x, y) , so that $z = f(x, y)$.

The function $z = f(x, y)$ defines a surface in \mathbb{R}^3 .

At any point (x, y) in the x - y plane, the direction in which one must travel in order to experience the greatest possible rate of increase in z at that point is the direction of the **gradient vector**,

$$\bar{\nabla}f = \frac{\partial f}{\partial x} \hat{\mathbf{i}} + \frac{\partial f}{\partial y} \hat{\mathbf{j}}$$

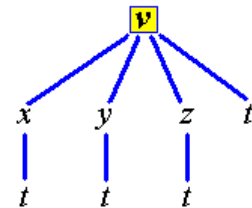
The magnitude of the gradient vector is that greatest possible rate of increase in z at that point. The gradient vector is not constant everywhere, unless the surface is a plane. (The symbol $\bar{\nabla}$ is usually pronounced “del”).

The concept of the gradient vector can be extended to functions of any number of

variables. If $u = f(x, y, z, t)$, then $\bar{\nabla}f = \left[\frac{\partial f}{\partial x} \quad \frac{\partial f}{\partial y} \quad \frac{\partial f}{\partial z} \quad \frac{\partial f}{\partial t} \right]^T$.

If \mathbf{v} is a function of position \mathbf{r} and time t , while position is in turn a function of time, then by the chain rule of differentiation,

$$\frac{d\bar{\mathbf{v}}}{dt} =$$



which is of use in the study of fluid dynamics.

The gradient operator can also be applied to vectors via the scalar (dot) and vector (cross) products:

The **divergence** of a vector field $\mathbf{F}(x, y, z)$ is

$$\operatorname{div} \bar{\mathbf{F}} = \bar{\nabla} \cdot \bar{\mathbf{F}} = \left[\frac{\partial}{\partial x} \quad \frac{\partial}{\partial y} \quad \frac{\partial}{\partial z} \right]^T \cdot [F_1 \ F_2 \ F_3]^T = \frac{\partial F_1}{\partial x} + \frac{\partial F_2}{\partial y} + \frac{\partial F_3}{\partial z}$$

A region free of sources and sinks will have zero divergence:
the total flux into any region is balanced by the total flux out from that region.

The **curl** of a vector field $\mathbf{F}(x, y, z)$ is

$$\operatorname{curl} \bar{\mathbf{F}} = \bar{\nabla} \times \bar{\mathbf{F}} = \begin{vmatrix} \hat{\mathbf{i}} & \frac{\partial}{\partial x} & F_1 \\ \hat{\mathbf{j}} & \frac{\partial}{\partial y} & F_2 \\ \hat{\mathbf{k}} & \frac{\partial}{\partial z} & F_3 \end{vmatrix} = \begin{bmatrix} \frac{\partial F_3}{\partial y} - \frac{\partial F_2}{\partial z} \\ \frac{\partial F_1}{\partial z} - \frac{\partial F_3}{\partial x} \\ \frac{\partial F_2}{\partial x} - \frac{\partial F_1}{\partial y} \end{bmatrix}$$

In an irrotational field, $\operatorname{curl} \bar{\mathbf{F}} = \bar{\mathbf{0}}$.

Whenever $\bar{\mathbf{F}} = \bar{\nabla} \phi$ for some twice differentiable potential function ϕ , $\operatorname{curl} \bar{\mathbf{F}} = \bar{\mathbf{0}}$
or

$$\operatorname{curl} (\operatorname{grad} \phi) \equiv \bar{\nabla} \times \bar{\nabla} \phi \equiv \bar{\mathbf{0}}$$

Proof:

$$\bar{\mathbf{F}} = \bar{\nabla} \phi = [F_1 \ F_2 \ F_3]^T = \left[\frac{\partial \phi}{\partial x} \quad \frac{\partial \phi}{\partial y} \quad \frac{\partial \phi}{\partial z} \right]^T$$

$$\Rightarrow \operatorname{curl} \bar{\nabla} \phi =$$

Among many identities involving the gradient operator is

$$\text{div}(\text{curl } \vec{\mathbf{F}}) \equiv \vec{\nabla} \cdot \vec{\nabla} \times \vec{\mathbf{F}} \equiv 0$$

for all twice-differentiable vector functions $\vec{\mathbf{F}}$

Proof:

$\text{div curl } \vec{\mathbf{F}} =$

The divergence of the gradient of a scalar function is the **Laplacian**:

$$\text{div}(\text{grad } f) \equiv \vec{\nabla} \cdot \vec{\nabla} f \equiv \nabla^2 f \equiv \frac{\partial^2 f}{\partial x^2} + \frac{\partial^2 f}{\partial y^2} + \frac{\partial^2 f}{\partial z^2}$$

for all twice-differentiable scalar functions f .

In orthogonal non-Cartesian coordinate systems, the expressions for the gradient operator are not as simple.

5.02 Differentiation in Orthogonal Curvilinear Coordinate Systems

For any orthogonal curvilinear coordinate system (u_1, u_2, u_3) in \mathbb{R}^3 , the unit tangent vectors along the curvilinear axes are $\hat{\mathbf{e}}_i = \hat{\mathbf{T}}_i = \frac{1}{h_i} \frac{\partial \bar{\mathbf{r}}}{\partial u_i}$,

where the scale factors $h_i = \left| \frac{\partial \bar{\mathbf{r}}}{\partial u_i} \right|$.

The displacement vector $\bar{\mathbf{r}}$ can then be written as $\bar{\mathbf{r}} = u_1 \hat{\mathbf{e}}_1 + u_2 \hat{\mathbf{e}}_2 + u_3 \hat{\mathbf{e}}_3$, where the unit vectors $\hat{\mathbf{e}}_i$ form an **orthonormal basis** for \mathbb{R}^3 .

$$\hat{\mathbf{e}}_i \cdot \hat{\mathbf{e}}_j = \delta_{ij} = \begin{cases} 0 & (i \neq j) \\ 1 & (i = j) \end{cases}$$

The differential displacement vector $d\mathbf{r}$ is (by the Chain Rule)

$$d\mathbf{r} = \frac{\partial \bar{\mathbf{r}}}{\partial u_1} du_1 + \frac{\partial \bar{\mathbf{r}}}{\partial u_2} du_2 + \frac{\partial \bar{\mathbf{r}}}{\partial u_3} du_3 = h_1 du_1 \hat{\mathbf{e}}_1 + h_2 du_2 \hat{\mathbf{e}}_2 + h_3 du_3 \hat{\mathbf{e}}_3$$

and the differential arc length ds is given by

$$ds^2 = d\bar{\mathbf{r}} \cdot d\bar{\mathbf{r}} = (h_1 du_1)^2 + (h_2 du_2)^2 + (h_3 du_3)^2$$

The element of volume dV is

$$dV = h_1 h_2 h_3 du_1 du_2 du_3 = \underbrace{\left| \frac{\partial(x, y, z)}{\partial(u_1, u_2, u_3)} \right|}_{\text{Jacobian}} du_1 du_2 du_3$$

$$= \begin{vmatrix} \frac{\partial x}{\partial u_1} & \frac{\partial y}{\partial u_1} & \frac{\partial z}{\partial u_1} \\ \frac{\partial x}{\partial u_2} & \frac{\partial y}{\partial u_2} & \frac{\partial z}{\partial u_2} \\ \frac{\partial x}{\partial u_3} & \frac{\partial y}{\partial u_3} & \frac{\partial z}{\partial u_3} \end{vmatrix} du_1 du_2 du_3$$

Example 5.02.1: Find the scale factor h_θ for the spherical polar coordinate system $(x, y, z) = (r \sin \theta \cos \phi, r \sin \theta \sin \phi, r \cos \theta)$:

5.03 Summary Table for the Gradient Operator

$$\text{Gradient operator} \quad \bar{\nabla} = \frac{\hat{\mathbf{e}}_1}{h_1} \frac{\partial}{\partial u_1} + \frac{\hat{\mathbf{e}}_2}{h_2} \frac{\partial}{\partial u_2} + \frac{\hat{\mathbf{e}}_3}{h_3} \frac{\partial}{\partial u_3}$$

$$\text{Gradient} \quad \bar{\nabla} V = \frac{\hat{\mathbf{e}}_1}{h_1} \frac{\partial V}{\partial u_1} + \frac{\hat{\mathbf{e}}_2}{h_2} \frac{\partial V}{\partial u_2} + \frac{\hat{\mathbf{e}}_3}{h_3} \frac{\partial V}{\partial u_3}$$

$$\text{Divergence} \quad \bar{\nabla} \cdot \bar{\mathbf{F}} = \frac{1}{h_1 h_2 h_3} \left(\frac{\partial (h_2 h_3 F_1)}{\partial u_1} + \frac{\partial (h_3 h_1 F_2)}{\partial u_2} + \frac{\partial (h_1 h_2 F_3)}{\partial u_3} \right)$$

$$\text{Curl} \quad \bar{\nabla} \times \bar{\mathbf{F}} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \hat{\mathbf{e}}_1 & \frac{\partial}{\partial u_1} & h_1 F_1 \\ h_2 \hat{\mathbf{e}}_2 & \frac{\partial}{\partial u_2} & h_2 F_2 \\ h_3 \hat{\mathbf{e}}_3 & \frac{\partial}{\partial u_3} & h_3 F_3 \end{vmatrix}$$

$$\text{Laplacian} \quad \nabla^2 V = \frac{1}{h_1 h_2 h_3} \left(\frac{\partial}{\partial u_1} \left(\frac{h_2 h_3}{h_1} \frac{\partial V}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left(\frac{h_3 h_1}{h_2} \frac{\partial V}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left(\frac{h_1 h_2}{h_3} \frac{\partial V}{\partial u_3} \right) \right)$$

Scale factors:

$$\text{Cartesian:} \quad h_x = h_y = h_z = 1 .$$

$$\text{Cylindrical polar:} \quad h_\rho = h_z = 1 , \quad h_\phi = \rho .$$

$$\text{Spherical polar:} \quad h_r = 1 , \quad h_\theta = r , \quad h_\phi = r \sin \theta .$$

Example 5.03.1: The Laplacian of V in spherical polars is

$$\nabla^2 V =$$

Example 5.03.2

A potential function $V(\vec{r})$ is spherically symmetric, (that is, its value depends only on the distance r from the origin), due solely to a point source at the origin. There are no other sources or sinks anywhere in \mathbb{R}^3 . Deduce the functional form of $V(\vec{r})$.

5.04 Derivatives of Basis Vectors

Cartesian: $\frac{d}{dt} \hat{\mathbf{i}} = \frac{d}{dt} \hat{\mathbf{j}} = \frac{d}{dt} \hat{\mathbf{k}} = \bar{\mathbf{0}}$ $\bar{\mathbf{r}} = x\hat{\mathbf{i}} + y\hat{\mathbf{j}} + z\hat{\mathbf{k}}$
 $\Rightarrow \bar{\mathbf{v}} = \dot{x}\hat{\mathbf{i}} + \dot{y}\hat{\mathbf{j}} + \dot{z}\hat{\mathbf{k}}$

Cylindrical Polar Coordinates:

$x = \rho \cos \phi, \quad y = \rho \sin \phi, \quad z = z$

$$\frac{d}{dt} \hat{\rho} = \frac{d\phi}{dt} \hat{\phi}$$

$$\frac{d}{dt} \hat{\phi} = -\frac{d\phi}{dt} \hat{\rho}$$

$$\frac{d}{dt} \hat{\mathbf{k}} = \bar{\mathbf{0}}$$

$$\mathbf{r} = \rho \hat{\rho} + z \hat{\mathbf{k}}$$

$$\Rightarrow \bar{\mathbf{v}} = \dot{\rho} \hat{\rho} + \rho \dot{\phi} \hat{\phi} + \dot{z} \hat{\mathbf{k}}$$

[radial and transverse components of $\bar{\mathbf{v}}$]

Spherical Polar Coordinates.

The “declination” angle θ is the angle between the positive z axis and the radius vector $\bar{\mathbf{r}}$. $0 \leq \theta \leq \pi$.

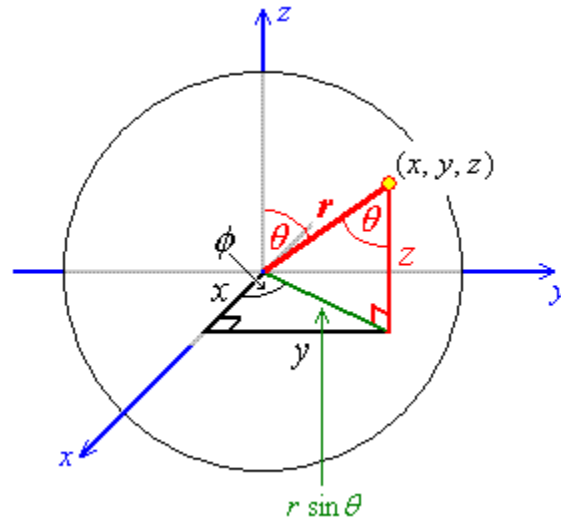
The “azimuth” angle ϕ is the angle on the x - y plane, measured anticlockwise from the positive x axis, of the shadow of the radius vector. $0 \leq \phi < 2\pi$.

$z = r \cos \theta.$

The shadow of the radius vector on the x - y plane has length $r \sin \theta$.

It then follows that

$x = r \sin \theta \cos \phi \quad \text{and} \quad y = r \sin \theta \sin \phi.$



$$\frac{d}{dt} \hat{\mathbf{r}} = \frac{d\theta}{dt} \hat{\theta} + \frac{d\phi}{dt} \sin \theta \hat{\phi}$$

$$\frac{d}{dt} \hat{\theta} = -\frac{d\theta}{dt} \hat{\mathbf{r}} + \frac{d\phi}{dt} \cos \theta \hat{\phi}$$

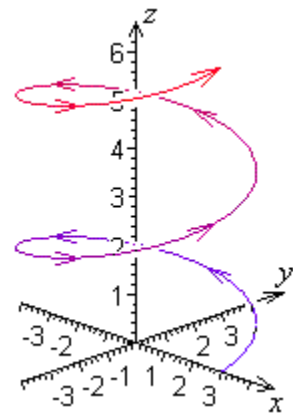
$$\frac{d}{dt} \hat{\phi} = -\frac{d\phi}{dt} (\sin \theta \hat{\mathbf{r}} + \cos \theta \hat{\theta})$$

$$\bar{\mathbf{r}} = r \hat{\mathbf{r}}$$

$$\Rightarrow \bar{\mathbf{v}} = \dot{r} \hat{\mathbf{r}} + r \dot{\theta} \hat{\theta} + r \dot{\phi} \sin \theta \hat{\phi}$$

Example 5.04.1

Find the velocity and acceleration in cylindrical polar coordinates for a particle travelling along the helix $x = 3 \cos 2t$, $y = 3 \sin 2t$, $z = t$.



Other examples are in the problem sets.