1.11 The Gamma Function

The gamma function $\Gamma(x)$ is a special function that will be needed in the solution of Bessel's ODE. $\Gamma(x)$ is a generalisation of the factorial function n! from positive integers to most real numbers. For any positive integer n, $n! = n \times (n-1) \times (n-2) \times ... \times 3 \times 2 \times 1$ (with 0! defined to be 1)

When x is a positive integer n, $\Gamma(n) = (n-1)!$

We know that $n! = n \times (n-1)!$

The gamma function has a similar recurrence relationship: $\Gamma(x+1) = x \cdot \Gamma(x)$

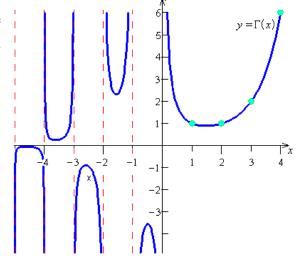
This allows $\Gamma(x)$ to be defined for non-integer negative x, using $\Gamma(x) = \frac{\Gamma(x+1)}{x}$ For example,

it can be shown that $\Gamma\left(\frac{1}{2}\right) = \sqrt{\pi}$

$$\Rightarrow \Gamma\left(-\frac{1}{2}\right) = \frac{\Gamma\left(+\frac{1}{2}\right)}{-\frac{1}{2}} = -2\sqrt{\pi} \Rightarrow \Gamma\left(-\frac{3}{2}\right) = \frac{\Gamma\left(-\frac{1}{2}\right)}{-\frac{3}{2}} = +\frac{4\sqrt{\pi}}{3}, \text{ etc.}$$

 $\Gamma(x)$ is infinite when x is a negative integer or zero. It is well defined for all other real numbers x.

In this graph of $y = \Gamma(x)$, values of the factorial function (at positive integer values of x) are highlighted.



There are several ways to define the gamma function, such as

$$\Gamma(x) = \int_0^\infty t^{x-1} e^{-t} dt \qquad (x > 0)$$

and

$$\Gamma(x) = \lim_{n \to \infty} \frac{n! n^x}{x(x+1)...(x+n)}$$

A related special function is the **beta function**:

$$B(m,n) = \int_0^1 t^{m-1} (1-t)^{n-1} dt = 2 \int_0^{\pi/2} \sin^{2m-1} \theta \cos^{2n-1} \theta d\theta = \frac{\Gamma(m)\Gamma(n)}{\Gamma(m+n)}$$

Among the many results involving the gamma function are:

For the closed region V in the first octant, bounded by the coordinate planes and the surface $\left(\frac{x}{a}\right)^{\alpha} + \left(\frac{y}{b}\right)^{\beta} + \left(\frac{z}{c}\right)^{\gamma} = 1$, with all constants positive,

$$I = \iiint_{V} x^{p-1} y^{q-1} z^{r-1} dx dy dz = \frac{a^{p} b^{q} c^{r}}{\alpha \beta \gamma} \cdot \frac{\Gamma\left(\frac{p}{\alpha}\right) \Gamma\left(\frac{q}{\beta}\right) \Gamma\left(\frac{r}{\gamma}\right)}{\Gamma\left(1 + \frac{p}{\alpha} + \frac{q}{\beta} + \frac{r}{\gamma}\right)}$$

For the closed area A in the first quadrant, bounded by the coordinate axes and the curve $\left(\frac{x}{a}\right)^{\alpha} + \left(\frac{y}{b}\right)^{\beta} = 1$, with all constants positive,

$$I = \iint_A x^{p-1} y^{q-1} dx dy = \frac{a^p b^q}{\alpha \beta} \cdot \frac{\Gamma\left(\frac{p}{\alpha}\right) \Gamma\left(\frac{q}{\beta}\right)}{\Gamma\left(1 + \frac{p}{\alpha} + \frac{q}{\beta}\right)}$$

Example 1.11.1

Establish the formula for the area enclosed by an ellipse.

The Cartesian equation of a standard ellipse is $\left(\frac{x}{a}\right)^2 + \left(\frac{y}{b}\right)^2 = 1$.

Set $\alpha = \beta = 2$ and p = q = 1, then

$$A = 4I = 4 \iint_{A} 1 \, dx \, dy = 4 \frac{a^{1}b^{1}}{2 \times 2} \cdot \frac{\Gamma(\frac{1}{2})\Gamma(\frac{1}{2})}{\Gamma(1 + \frac{1}{2} + \frac{1}{2})} = ab \cdot \frac{\left(\Gamma(\frac{1}{2})\right)^{2}}{\Gamma(2)} = ab \frac{\left(\sqrt{\pi}\right)^{2}}{1!} = \pi ab$$

1.12 Bessel and Legendre ODEs

Frobenius Series Solution of an ODE

If the ODE

$$P(x)y'' + Q(x)y' + R(x)y = F(x)$$

is such that $P(x_0) = 0$, but $(x - x_0) \frac{Q(x)}{P(x)}$, $(x - x_0)^2 \frac{R(x)}{P(x)}$ and $\frac{F(x)}{P(x)}$ are all analytic at x_0 , then $x = x_0$ is a **regular singular point** of the ODE.

A Frobenius series solution of the ODE about $x = x_0$ exists:

$$y(x) = \sum_{n=0}^{\infty} c_n (x - x_0)^{n+r}$$

for some real number(s) r and for some set of values $\{c_n\}$.

Example 1.12.1

Find a solution of Bessel's ordinary differential equation of order v, ($v \ge 0$),

$$x^2y'' + xy' + (x^2 - v^2)y = 0$$

 $P(0) = 0 \implies x_0 = 0$ is a singular point.

$$(x-x_0)\frac{Q(x)}{P(x)} = x\frac{x}{x^2} = 1$$
, $(x-x_0)^2 \frac{R(x)}{P(x)} = x^2 \frac{(x^2-v^2)}{x^2} = x^2-v^2$
and $\frac{F(x)}{P(x)} = 0$

Therefore $x_0 = 0$ is a regular singular point of Bessel's equation.

Substitute the Frobenius series $y(x) = \sum_{n=0}^{\infty} c_n x^{n+r}$ into the ODE:

$$\sum_{n=0}^{\infty} (n+r)(n+r-1)c_n x^{n+r-2+2} + \sum_{n=0}^{\infty} (n+r)c_n x^{n+r-1+1} + \sum_{n=0}^{\infty} c_n x^{n+r+2} - \sum_{n=0}^{\infty} v^2 c_n x^{n+r} = 0$$

Example 1.12.1 (continued)

Adjust the index on the third summation so that the exponents of *x* match:

$$\sum_{n=0}^{\infty} c_n \Big[(n+r)(n+r-1+1) - v^2 \Big] x^{n+r} + \sum_{n=2}^{\infty} c_{n-2} x^{n+r} = 0$$

The summations can be combined for n = 2 onwards:

$$\left(r^2 - v^2\right)c_0x^r + \left[\left(r+1\right)^2 - v^2\right]c_1x^{r+1} + \sum_{n=2}^{\infty} \left[\left(\left(n+r\right)^2 - v^2\right)c_n + c_{n-2}\right]x^{n+r} = 0$$

Setting the coefficient of x^r (the lowest exponent present) to zero generates the **indicial** equation $r^2 - v^2 = 0 \implies r = \pm v$.

Examining the positive root, the series now becomes

$$0 + \left[\left(v + 1 \right)^2 - v^2 \right] c_1 x^{r+1} + \sum_{n=2}^{\infty} \left[\left(\left(n + v \right)^2 - v^2 \right) c_n + c_{n-2} \right] x^{n+r} = 0$$

$$\Rightarrow (2v+1) c_1 x^{r+1} + \sum_{n=2}^{\infty} \left[\left(2nv + n^2 \right) c_n + c_{n-2} \right] x^{n+r} = 0$$
But $v \ge 0 \Rightarrow 2v + 1 \ne 0 \Rightarrow c_1 = 0$
For $n \ge 2$, $c_n = \frac{-c_{n-2}}{n(n+2v)}$

It then follows that this series must be even: $0 = c_1 = c_3 = c_5 = \dots$ or $c_{2k-1} = 0 \quad \forall k \in \mathbb{N}$ For the even order terms, replace the index n by the even index 2k (where k is any natural number) and pursue the recurrence relation down to c_0 :

$$c_{2k} = \frac{-c_{2k-2}}{2k(2k+2\nu)} = \frac{(-1)}{2^2k(k+\nu)}c_{2(k-1)} = \frac{(-1)}{2^2k(k+\nu)} \cdot \frac{(-1)}{2^2(k-1)(k-1+\nu)}c_{2(k-2)}$$

$$= \frac{(-1)}{2^2k(k+\nu)} \cdot \frac{(-1)}{2^2(k-1)(k-1+\nu)} \cdot \frac{(-1)}{2^2(k-2)(k-2+\nu)}c_{2(k-3)} = \dots$$

$$= \frac{(-1)^k}{2^{2k}k(k-1)(k-2)\dots(k-[k-1])\cdot(k+\nu)(k-1+\nu)(k-2+\nu)\dots(k-[k-1]+\nu)}c_{2(k-k)}$$

$$\Rightarrow c_{2k} = \frac{(-1)^k}{2^{2k}k!(\nu+k)(\nu+k-1)\dots(\nu+1)}c_0$$

$$\Rightarrow c_{2k} = \frac{(-1)^k}{2^{2k}k!(\nu+k)(\nu+k-1)\dots(\nu+1)} \cdot \frac{\Gamma(\nu+1)}{\Gamma(\nu+1)}c_0$$

$$\Rightarrow c_{2k} = \frac{(-1)^k}{2^{2k}k!(\nu+k)(\nu+k-1)\dots(\nu+1)} \cdot \frac{\Gamma(\nu+1)}{\Gamma(\nu+1)}c_0$$

$$\Rightarrow c_{2k} = \frac{(-1)^k}{2^{2k}k!(\nu+k)(\nu+k-1)\dots(\nu+1)} \cdot \frac{\Gamma(\nu+1)}{\Gamma(\nu+1)}c_0$$

Example 1.12.1 (continued)

One Frobenius solution of Bessel's equation of order v is therefore

$$y(x) = c_0 \Gamma(\nu+1) \sum_{k=0}^{\infty} \frac{(-1)^k}{2^{2k} k! \Gamma(\nu+k+1)} x^{2k+\nu} = c_0 \Gamma(\nu+1) J_{\nu}(x)$$

where $J_{\nu}(x)$ is the **Bessel function of the first kind of order** ν .

It turns out that the Frobenius series found by setting r = -v generates a second linearly independent solution $J_{-v}(x)$ of the Bessel equation only if v is not an integer.

The Bessel ODE in standard form,

$$x^2y'' + xy' + (x^2 - v^2)y = 0$$

has the general solution

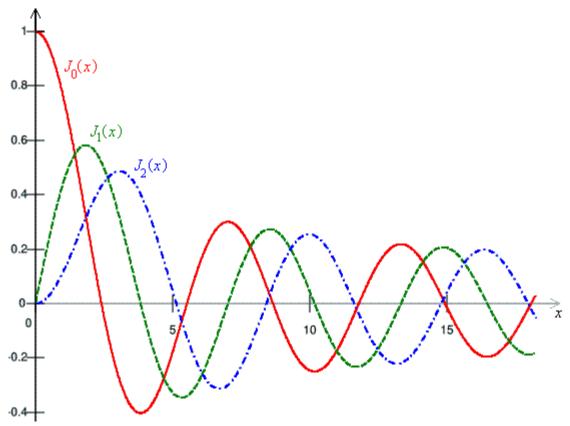
$$y(x) = A J_{\nu}(x) + B Y_{\nu}(x)$$

unless ν is not an integer, in which case $Y_{\nu}(x)$ can be replaced by $J_{-\nu}(x)$.

 $Y_{\nu}(x)$ is the Bessel function of the second kind.

When ν is an integer, $J_{-\nu}(x) = (-1)^{\nu} J_{\nu}(x)$.

Graphs of Bessel functions of the first kind, for v = 0, 1, 2:



The series expression for the Bessel function of the first kind is

$$J_{\nu}(x) = \sum_{k=0}^{\infty} \frac{\left(-1\right)^{k}}{k! \Gamma(k+\nu+1)} \left(\frac{x}{2}\right)^{2k+\nu}$$

This function has a simpler form when v is an odd half-integer. For example,

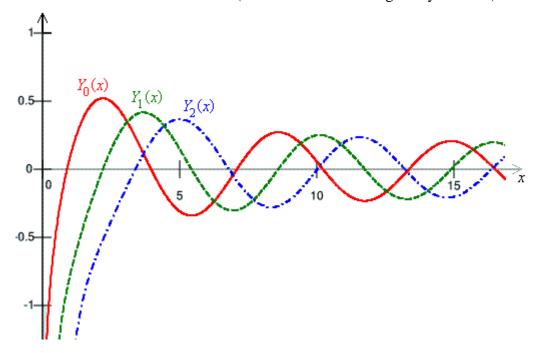
$$J_{1/2}(x) = \sqrt{\frac{2}{\pi x}} \sin x, \qquad J_{-1/2}(x) = \sqrt{\frac{2}{\pi x}} \cos x$$

The Bessel function of the second kind is

$$Y_{\nu}(x) = \frac{J_{\nu}(x)\cos(\nu\pi) - J_{-\nu}(x)}{\sin(\nu\pi)}$$

 $Y_{\nu}(x)$ is unbounded as $x \to 0$: $\lim_{x \to 0^+} Y_{\nu}(x) = -\infty$

Bessel functions of the second kind (all of which have a singularity at x = 0):



Bessel functions arise frequently in situations where cylindrical or spherical polar coordinates are used.

A generalised Bessel ODE is

$$x^{2} \frac{d^{2}y}{dx^{2}} + (1 - 2a)x \frac{dy}{dx} + (b^{2}c^{2}x^{2c} + (a^{2} - c^{2}v^{2}))y = 0$$

whose general solution is

$$y(x) = x^{a} \left(A J_{\nu} \left(b x^{c} \right) + B Y_{\nu} \left(b x^{c} \right) \right)$$

For a generalised Bessel ODE with $a \ge 0$, whenever the solution must remain bounded as $x \to 0$, the general solution simplifies to $y(x) = Ax^a J_v(bx^c)$.

Example 1.12.2

Find a Maclaurin series solution to Legendre's ODE

$$(1-x^2)\frac{d^2y}{dx^2} - 2x\frac{dy}{dx} + p(p+1)y = 0$$

in the case when p is a non-negative integer.

$$P(x)=1-x^2 \implies P(0)=1\neq 0 \implies x=0$$
 is a regular point of the ODE.

Let the general solution be $y(x) = \sum_{n=0}^{\infty} a_n x^n$.

Then
$$y'(x) = \sum_{n=1}^{\infty} n a_n x^{n-1}$$
 and $y''(x) = \sum_{n=2}^{\infty} n(n-1) a_n x^{n-2}$.

Substitute into the ODE:

$$\sum_{n=0}^{\infty} n(n-1)a_n \left(x^{n-2} - x^n\right) - \sum_{n=0}^{\infty} 2n a_n x^n + \sum_{n=0}^{\infty} p(p+1)a_n x^n = 0$$

$$\Rightarrow \sum_{n=0}^{\infty} n(n-1)a_n x^{n-2} + \sum_{n=0}^{\infty} (-n(n-1) - 2n + p(p+1))a_n x^n = 0$$

But the first two terms (n = 0 and n = 1) of the first series are both zero.

$$\sum_{n=0}^{\infty} n(n-1)a_n x^{n-2} = 0 + 0 + \sum_{n=2}^{\infty} n(n-1)a_n x^{n-2} = \sum_{n=0}^{\infty} (n+2)(n+1)a_{n+2} x^n$$

after a shift in indices. Returning to the full ODE,

$$\sum_{n=0}^{\infty} (n+2)(n+1)a_{n+2}x^{n} + \sum_{n=0}^{\infty} (-n^{2}+n-2n + p^{2}+p)a_{n}x^{n} = 0$$

$$\Rightarrow \sum_{n=0}^{\infty} \left((n+2)(n+1)a_{n+2} + \left(-(n^2 - p^2) - (n-p) \right) a_n \right) x^n = 0$$

Matching coefficients of x^n , $(n \ge 0)$:

$$(n+2)(n+1)a_{n+2} = (n^2 - p^2 + n - p)a_n \implies a_{n+2} = \frac{(n-p)(n+p+1)}{(n+2)(n+1)}a_n$$

Shifting indices, $a_n = \frac{(n+p-1)(n-p-2)}{n(n-1)}a_{n-2}$ $(n \ge 2)$, with a_0 and a_1 arbitrary.

Rearranging slightly,
$$a_n = \frac{-(p-(n-2))(p+(n-1))}{n(n-1)}a_{n-2} \quad (n \ge 2)$$

Example 1.12.2 (continued)

$$\Rightarrow a_2 = \frac{-p(p+1)}{2\times 1} a_0, \quad a_3 = \frac{-(p-1)(p+2)}{3\times 2} a_1,$$

$$a_4 = \frac{-(p-2)(p+3)}{4\times 3} a_2 = \frac{+p(p+1)(p-2)(p+3)}{4\times 3\times 2\times 1} a_0,$$

$$a_5 = \frac{-(p-3)(p+4)}{5\times 4} a_3 = \frac{+(p-1)(p+2)(p-3)(p+4)}{5\times 4\times 3\times 2\times 1} a_1,$$

$$a_6 = \frac{-(p-4)(p+5)}{6\times 5} a_4 = \frac{-p(p+1)(p-2)(p+3)(p-4)(p+5)}{6\times 5\times 4\times 3\times 2\times 1} a_0, \text{ etc.}$$

It then follows that the general solution to Legendre's ODE is

$$y_{p}(x) = a_{0} \left(1 - \frac{p(p+1)}{2!} x^{2} + \frac{p(p+1)(p-2)(p+3)}{4!} x^{4} - \frac{p(p+1)(p-2)(p+3)(p-4)(p+5)}{6!} x^{6} + \dots \right)$$

$$+ a_{1} \left(x - \frac{(p-1)(p+2)}{3!} x^{3} + \frac{(p-1)(p+2)(p-3)(p+4)}{5!} x^{5} - \frac{(p-1)(p+2)(p-3)(p+4)(p-5)(p+6)}{7!} x^{7} + \dots \right)$$

where a_0 and a_1 are arbitrary constants. This series converges on [-1, 1].

These solutions $y(x) = \sum_{n=0}^{\infty} a_n x^n$ are Legendre functions of order p.

If p is a non-negative integer then

$$a_{p+2} = \frac{(p-p)(p+p+1)}{(p+2)(p+1)} a_p = 0 \implies a_{p+4} = 0 \implies \dots$$

$$\Rightarrow a_{p+2k} = 0 \quad \forall k \in \mathbb{N}$$

If we set $a_1 = 0$ when p is even, then the series solution terminates as a pth order polynomial (and therefore converges for all x).

If we set $a_0 = 0$ when p is odd, then the series solution terminates as a pth order polynomial (and therefore converges for all x).

With suitable choices of a_0 and a_1 , so that $P_n(1) = 1$,

we have the set of **Legendre polynomials**:

$$P_{0}(x) = 1, P_{1}(x) = x, P_{2}(x) = \frac{1}{2}(3x^{2} - 1),$$

$$P_{3}(x) = \frac{1}{2}(5x^{3} - 3x), P_{4}(x) = \frac{1}{8}(35x^{4} - 30x^{2} + 3), P_{5}(x) = \frac{1}{8}(63x^{5} - 70x^{3} + 15x),$$

$$P_{6}(x) = \frac{1}{16}(231x^{6} - 315x^{4} + 105x^{2} - 5), \text{ etc.}$$

Each $P_n(x)$ is a solution of Legendre's ODE with p = n.

Rodrigues' formula generates all of the Legendre polynomials:

$$P_n(x) = \frac{1}{2^n n!} \frac{d^n}{dx^n} \left(\left(x^2 - 1 \right)^n \right)$$

Among the properties of Legendre polynomials is their orthogonality on [-1, 1]:

$$\int_{-1}^{+1} P_m(x) P_n(x) dx = \begin{cases} 0 & (m \neq n) \\ \frac{2}{2n+1} & (m=n) \end{cases}$$

[Space for any additional notes]