

Example 8.04.4

Find the complete solution to

$$\frac{\partial^2 u}{\partial x^2} + 2 \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial^2 u}{\partial y^2} = 0,$$

$$u = 0 \text{ on } x = 0,$$

$$u = x^2 \text{ on } y = 1.$$

$$A = 1, \quad B = 2, \quad C = 1 \quad \Rightarrow \quad D = 4 - 4 \times 1 = 0$$

Therefore the PDE is **parabolic** everywhere.

$$\lambda = \frac{-2 \pm \sqrt{0}}{2} = -1 \text{ or } -1$$

The complementary function (and general solution) is

$$u(x, y) = f(y-x) + h(x, y)g(y-x)$$

where $h(x, y)$ is any convenient non-trivial linear function of (x, y) except a multiple of $(y-x)$. Choosing, arbitrarily, $h(x, y) = x$,

$$u(x, y) = f(y-x) + xg(y-x)$$

Imposing the boundary conditions:

$$u(0, y) = 0 \quad \Rightarrow \quad f(y) + 0 = 0$$

Therefore the function f is identically zero, for any argument including $(y-x)$.

We now have $u(x, y) = xg(y-x)$.

$$u(x, 1) = x^2 \quad \Rightarrow \quad xg(1-x) = x^2 \quad \Rightarrow \quad g(1-x) = x \quad \Rightarrow \quad g(x) = 1-x$$

Therefore

$$u(x, y) = xg(y-x) = x(1-(y-x))$$

The complete solution is

$$\boxed{u(x, y) = x(x-y+1)}$$

Example 8.04.4 - Alternative Treatment of the Complementary Function

Find the complete solution to

$$\frac{\partial^2 u}{\partial x^2} + 2 \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial^2 u}{\partial y^2} = 0,$$

$$u = 0 \text{ on } x = 0,$$

$$u = x^2 \text{ on } y = 1.$$

$$A = 1, \quad B = 2, \quad C = 1 \quad \Rightarrow \quad D = 4 - 4 \times 1 = 0$$

$$\lambda = \frac{-2 \pm \sqrt{0}}{2} = -1 \text{ or } -1$$

The complementary function (and general solution) is

$$u(x, y) = f(y-x) + h(x, y)g(y-x)$$

where $h(x, y)$ is any convenient non-trivial linear function of (x, y) except a multiple of $(y-x)$. The most general choice possible is $h(x, y) = ax + by$, with the restriction $a \neq \lambda b$.

$$u(x, y) = f(y-x) + (ax+by) \cdot g(y-x)$$

$$\Rightarrow u(0, y) = f(y) + by \cdot g(y) = 0 \quad \Rightarrow \quad f(y) = -by \cdot g(y) \quad \text{(A)}$$

[note: choosing $b=0 \rightarrow f(y)=0$, as happened on the previous page] and

$$u(x, 1) = f(1-x) + (ax+b) \cdot g(1-x) = x^2$$

Using (A),

$$u(x, 1) = -b(1-x) \cdot g(1-x) + (ax+b) \cdot g(1-x) = x^2$$

$$\Rightarrow (-\cancel{b} + bx + ax + \cancel{b}) \cdot g(1-x) = x^2 \quad \Rightarrow \quad (a+b)x \cdot g(1-x) = x^2$$

$$\Rightarrow g(1-x) = \frac{x}{a+b} \quad (\text{note that the restriction on } h(x, y) \text{ ensures that } a+b \neq 0).$$

$$\Rightarrow g(y) = \frac{1-y}{a+b} \quad \text{(B)}$$

$$\text{(A)} \Rightarrow f(y) = -by \cdot \frac{1-y}{a+b}$$

The complete solution becomes

$$u(x, y) = -b(y-x) \cdot \frac{1-(y-x)}{a+b} + (ax+by) \frac{1-(y-x)}{a+b}$$

$$= (-\cancel{b}y + bx + ax + \cancel{b}y) \frac{1-(y-x)}{a+b} = (\cancel{a+b})x \frac{x-y+1}{\cancel{a+b}} \Rightarrow$$

$$u(x, y) = x(x-y+1)$$

Example 8.04.4 Extension (continued)

It is easy to confirm that this solution is correct:

$$u(x, y) = x(x - y + 1) \Rightarrow$$

$$u(0, y) = 0(0 - y + 1) = 0 \quad \forall y \quad \checkmark$$

and

$$u(x, 1) = x(x - 1 + 1) = x^2 \quad \forall x \quad \checkmark$$

and

$$u(x, y) = x^2 - xy + x \Rightarrow \frac{\partial u}{\partial x} = 2x - y \quad \text{and} \quad \frac{\partial u}{\partial y} = -x$$

$$\Rightarrow \frac{\partial^2 u}{\partial x^2} = 2, \quad \frac{\partial^2 u}{\partial x \partial y} = -1 \quad \text{and} \quad \frac{\partial^2 u}{\partial y^2} = 0$$

$$\Rightarrow \frac{\partial^2 u}{\partial x^2} + 2 \frac{\partial^2 u}{\partial x \partial y} + \frac{\partial^2 u}{\partial y^2} = 2 - 2 + 0 = 0 \quad \forall (x, y) \quad \checkmark$$

Two-dimensional Laplace Equation

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

$$A = C = 1, \quad B = 0 \quad \Rightarrow \quad D = 0 - 4 < 0$$

This PDE is **elliptic** everywhere.

$$\lambda = \frac{0 \pm \sqrt{-4}}{2} = \pm j$$

The general solution is

$$u(x, y) = f(y - jx) + g(y + jx)$$

where f and g are any twice-differentiable functions.

A function $f(x, y)$ is **harmonic** if and only if $\nabla^2 f = 0$ everywhere inside a domain Ω .

Example 8.04.5

Is $u = e^x \sin y$ harmonic on \mathbb{R}^2 ?

$$\frac{\partial u}{\partial x} = e^x \sin y \quad \text{and} \quad \frac{\partial u}{\partial y} = e^x \cos y$$

$$\Rightarrow \frac{\partial^2 u}{\partial x^2} = e^x \sin y \quad \text{and} \quad \frac{\partial^2 u}{\partial y^2} = -e^x \sin y$$

$$\Rightarrow \nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = e^x \sin y - e^x \sin y = 0 \quad \forall (x, y)$$

Therefore **yes**, $u = e^x \sin y$ is harmonic on \mathbb{R}^2 .

Example 8.04.6

Find the complete solution $u(x, y)$ to the partial differential equation $\nabla^2 u = 0$, given the additional information

$$u(0, y) = y^3 \quad \text{and} \quad \left. \frac{\partial u}{\partial x} \right|_{x=0} = 0$$

The PDE is

$$\nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

(which means that the solution $u(x, y)$ is an harmonic function).

$$\Rightarrow A = C = 1, \quad B = 0 \quad \Rightarrow D = B^2 - 4AC = -4 < 0$$

The PDE is **elliptic** everywhere.

$$\text{A.E.: } \lambda^2 + 1 = 0 \quad \Rightarrow \lambda = \pm j$$

$$\text{C.F.: } u_c(x, y) = f(y - jx) + g(y + jx)$$

The PDE is homogeneous \Rightarrow

$$\text{P.S.: } u_p(x, y) = 0$$

$$\text{G.S.: } u(x, y) = f(y - jx) + g(y + jx)$$

$$\Rightarrow u_x(x, y) = -j f'(y - jx) + j g'(y + jx)$$

Using the additional information,

$$u(0, y) = f(y) + g(y) = y^3 \quad \Rightarrow g(y) = y^3 - f(y) \quad \Rightarrow g'(y) = 3y^2 - f'(y)$$

and

$$u_x(0, y) = 0 = -j f'(y) + j g'(y) = j(-f'(y) + 3y^2 - f'(y))$$

$$\Rightarrow 2f'(y) = 3y^2 \quad \Rightarrow 2f(y) = y^3 \quad \Rightarrow f(y) = \frac{1}{2}y^3$$

$$\Rightarrow g(y) = y^3 - \frac{1}{2}y^3 = \frac{1}{2}y^3$$

Therefore the complete solution is

$$u(x, y) = \frac{1}{2}(y - jx)^3 + \frac{1}{2}(y + jx)^3$$

$$= \frac{1}{2} \left(y^3 - 3(jx)y^2 + 3(jx)^2 y - (jx)^3 + y^3 + 3(jx)y^2 + 3(jx)^2 y + (jx)^3 \right)$$

$$= y^3 + 3j^2 x^2 y \quad \Rightarrow$$

$$\boxed{u(x, y) = y^3 - 3x^2 y}$$

Note that the solution is completely real, even though the eigenvalues are not real.

8.05 The Maximum-Minimum Principle

Let Ω be some finite domain on which a function $u(x, y)$ and its second derivatives are defined. Let $\bar{\Omega}$ be the union of the domain with its boundary.

Let m and M be the minimum and maximum values respectively of u on the boundary of the domain.

If $\nabla^2 u \geq 0$ in Ω , then u is **subharmonic** and

$$u(\bar{\mathbf{r}}) < M \quad \text{or} \quad u(\bar{\mathbf{r}}) \equiv M \quad \forall \bar{\mathbf{r}} \text{ in } \Omega$$

If $\nabla^2 u \leq 0$ in Ω , then u is **superharmonic** and

$$u(\bar{\mathbf{r}}) > m \quad \text{or} \quad u(\bar{\mathbf{r}}) \equiv m \quad \forall \bar{\mathbf{r}} \text{ in } \Omega$$

If $\nabla^2 u = 0$ in Ω , then u is **harmonic** (both subharmonic and superharmonic) and u is either constant on $\bar{\Omega}$ or $m < u < M$ everywhere on Ω .

Example 8.05.1

$\nabla^2 u = 0$ in $\Omega : x^2 + y^2 < 1$ and $u(x, y) = 1$ on $C : x^2 + y^2 = 1$.
Find $u(x, y)$ on Ω .

$$u \text{ is harmonic on } \Omega \Rightarrow \min_C(u) \leq \left(\begin{array}{c} u(x, y) \\ \text{on } \Omega \end{array} \right) \leq \max_C(u)$$

But $\min_C(u) = \max_C(u) = 1$

Therefore $u(x, y) = 1$ everywhere in Ω .

Example 8.05.2

$\nabla^2 u = 0$ in the square domain $\Omega : -2 < x < +2, -2 < y < +2$.

On the boundary C , on the left and right edges ($x = \pm 2$), $u(x, y) = 4 - y^2$,
while on the top and bottom edges ($y = \pm 2$), $u(x, y) = x^2 - 4$.

Find bounds on the value of $u(x, y)$ inside the domain Ω .

For $-2 \leq y \leq +2$, $0 \leq 4 - y^2 \leq 4$.

For $-2 \leq x \leq +2$, $-4 \leq x^2 - 4 \leq 0$.

Therefore, on the boundary C of the domain Ω , $-4 \leq u(x, y) \leq +4$ so that
 $m = -4$ and $M = +4$.

$u(x, y)$ is harmonic (because $\nabla^2 u = 0$).

Therefore, everywhere in Ω ,

$$-4 < u(x, y) < +4$$

Note:

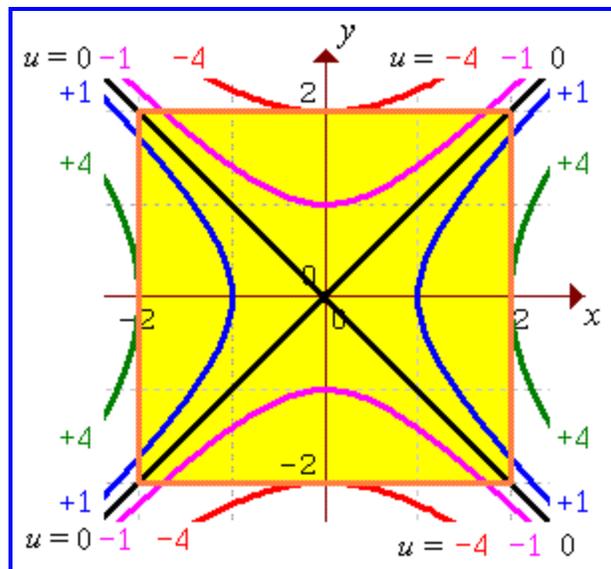
$u(x, y) = x^2 - y^2$ is consistent with the boundary condition and

$$\frac{\partial u}{\partial x} = 2x - 0, \quad \frac{\partial u}{\partial y} = 0 - 2y \Rightarrow \frac{\partial^2 u}{\partial x^2} = 2, \quad \frac{\partial^2 u}{\partial y^2} = -2 = -\frac{\partial^2 u}{\partial x^2}$$

$$\Rightarrow \nabla^2 u = \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0$$

Contours of constant values of u are hyperbolas.

A contour map illustrates that $-4 < u(x, y) < +4$ within the domain is indeed true.



8.06 The Heat Equation

For a material of constant density ρ , constant specific heat μ and constant thermal conductivity K , the partial differential equation governing the temperature u at any location (x, y, z) and any time t is

$$\frac{\partial u}{\partial t} = k \nabla^2 u, \quad \text{where } k = \frac{K}{\mu\rho}$$

Example 8.06.1

Heat is conducted along a thin homogeneous bar extending from $x = 0$ to $x = L$. There is no heat loss from the sides of the bar. The two ends of the bar are maintained at temperatures T_1 (at $x = 0$) and T_2 (at $x = L$). The initial temperature throughout the bar at the cross-section x is $f(x)$.

Find the temperature at any point in the bar at any subsequent time.

The partial differential equation governing the temperature $u(x, t)$ in the bar is

$$\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$$

together with the boundary conditions

$$u(0, t) = T_1 \quad \text{and} \quad u(L, t) = T_2$$

and the initial condition

$$u(x, 0) = f(x)$$

[Note that if an end of the bar is insulated, instead of being maintained at a constant temperature, then the boundary condition changes to $\frac{\partial u}{\partial x}(0, t) = 0$ or $\frac{\partial u}{\partial x}(L, t) = 0$.]

Attempt a solution by the method of separation of variables.

$$u(x, t) = X(x) T(t)$$

$$\Rightarrow XT' = kX''T \quad \Rightarrow \quad \frac{T'}{T} = k \frac{X''}{X} = c$$

Again, when a function of t only equals a function of x only, both functions must equal the same absolute constant. Unfortunately, the two boundary conditions cannot both be satisfied unless $T_1 = T_2 = 0$. Therefore we need to treat this more general case as a perturbation of the simpler ($T_1 = T_2 = 0$) case.

Example 8.06.1 (continued)

Let $u(x, t) = v(x, t) + g(x)$

Substitute this into the PDE:

$$\frac{\partial}{\partial t}(v(x, t) + g(x)) = k \frac{\partial^2}{\partial x^2}(v(x, t) + g(x)) \Rightarrow \frac{\partial v}{\partial t} = k \left(\frac{\partial^2 v}{\partial x^2} + g''(x) \right)$$

This is the standard heat PDE for v if we choose g such that $g''(x) = 0$.
 $g(x)$ must therefore be a linear function of x .

We want the perturbation function $g(x)$ to be such that

$$u(0, t) = T_1 \quad \text{and} \quad u(L, t) = T_2$$

and

$$v(0, t) = v(L, t) = 0$$

Therefore $g(x)$ must be the linear function for which $g(0) = T_1$ and $g(L) = T_2$.

It follows that

$$g(x) = \left(\frac{T_2 - T_1}{L} \right) x + T_1$$

and we now have the simpler problem

$$\frac{\partial v}{\partial t} = k \frac{\partial^2 v}{\partial x^2}$$

together with the boundary conditions

$$v(0, t) = v(L, t) = 0$$

and the initial condition

$$v(x, 0) = f(x) - g(x)$$

Now try separation of variables on $v(x, t)$:

$$v(x, t) = X(x) T(t)$$

$$\Rightarrow X T' = k X'' T \quad \Rightarrow \frac{1}{k} \frac{T'}{T} = \frac{X''}{X} = c$$

$$\text{But } v(0, t) = v(L, t) = 0 \Rightarrow X(0) = X(L) = 0$$

This requires c to be a negative constant, say $-\lambda^2$.

The solution is very similar to that for the wave equation on a finite string with fixed ends

(section 8.02). The eigenvalues are $\lambda = \frac{n\pi}{L}$ and the corresponding eigenfunctions are

any non-zero constant multiples of

$$X_n(x) = \sin\left(\frac{n\pi x}{L}\right)$$

Example 8.06.1 (continued)

The ODE for $T(t)$ becomes

$$T' + \left(\frac{n\pi}{L}\right)^2 kT = 0$$

whose general solution is

$$T_n(t) = c_n e^{-n^2\pi^2 kt/L^2}$$

Therefore

$$v_n(x,t) = X_n(x)T_n(t) = c_n \sin\left(\frac{n\pi x}{L}\right) \exp\left(-\frac{n^2\pi^2 kt}{L^2}\right)$$

If the initial temperature distribution $f(x) - g(x)$ is a simple multiple of $\sin\left(\frac{n\pi x}{L}\right)$ for

some integer n , then the solution for v is just $v(x,t) = c_n \sin\left(\frac{n\pi x}{L}\right) \exp\left(-\frac{n^2\pi^2 kt}{L^2}\right)$.

Otherwise, we must attempt a superposition of solutions.

$$v(x,t) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{L}\right) \exp\left(-\frac{n^2\pi^2 kt}{L^2}\right)$$

such that $v(x,0) = \sum_{n=1}^{\infty} c_n \sin\left(\frac{n\pi x}{L}\right) = f(x) - g(x)$.

The Fourier sine series coefficients are $c_n = \frac{2}{L} \int_0^L (f(z) - g(z)) \sin\left(\frac{n\pi z}{L}\right) dz$

so that the complete solution for $v(x,t)$ is

$$v(x,t) = \frac{2}{L} \sum_{n=1}^{\infty} \left(\int_0^L \left(f(z) - \frac{T_2 - T_1}{L} z - T_1 \right) \sin\left(\frac{n\pi z}{L}\right) dz \right) \sin\left(\frac{n\pi x}{L}\right) \exp\left(-\frac{n^2\pi^2 kt}{L^2}\right)$$

and the complete solution for $u(x,t)$ is

$$u(x,t) = v(x,t) + \left(\frac{T_2 - T_1}{L}\right)x + T_1$$

Note how this solution can be partitioned into a transient part $v(x,t)$ (which decays to zero as t increases) and a steady-state part $g(x)$ which is the limiting value that the temperature distribution approaches.

Example 8.06.1 (continued)

As a specific example, let $k = 9$, $T_1 = 100$, $T_2 = 200$, $L = 2$ and

$f(x) = 145x^2 - 240x + 100$, (for which $f(0) = 100$, $f(2) = 200$ and $f(x) > 0 \forall x$).

Then $g(x) = \frac{200-100}{2}x + 100 = 50x + 100$

The Fourier sine series coefficients are

$$c_n = \int_0^2 \left((145z^2 - 240z + 100) - (50z + 100) \right) \sin\left(\frac{n\pi z}{2}\right) dz$$

$$\Rightarrow c_n = 145 \int_0^2 (z^2 - 2z) \sin\left(\frac{n\pi z}{2}\right) dz$$

After an integration by parts (details omitted here),

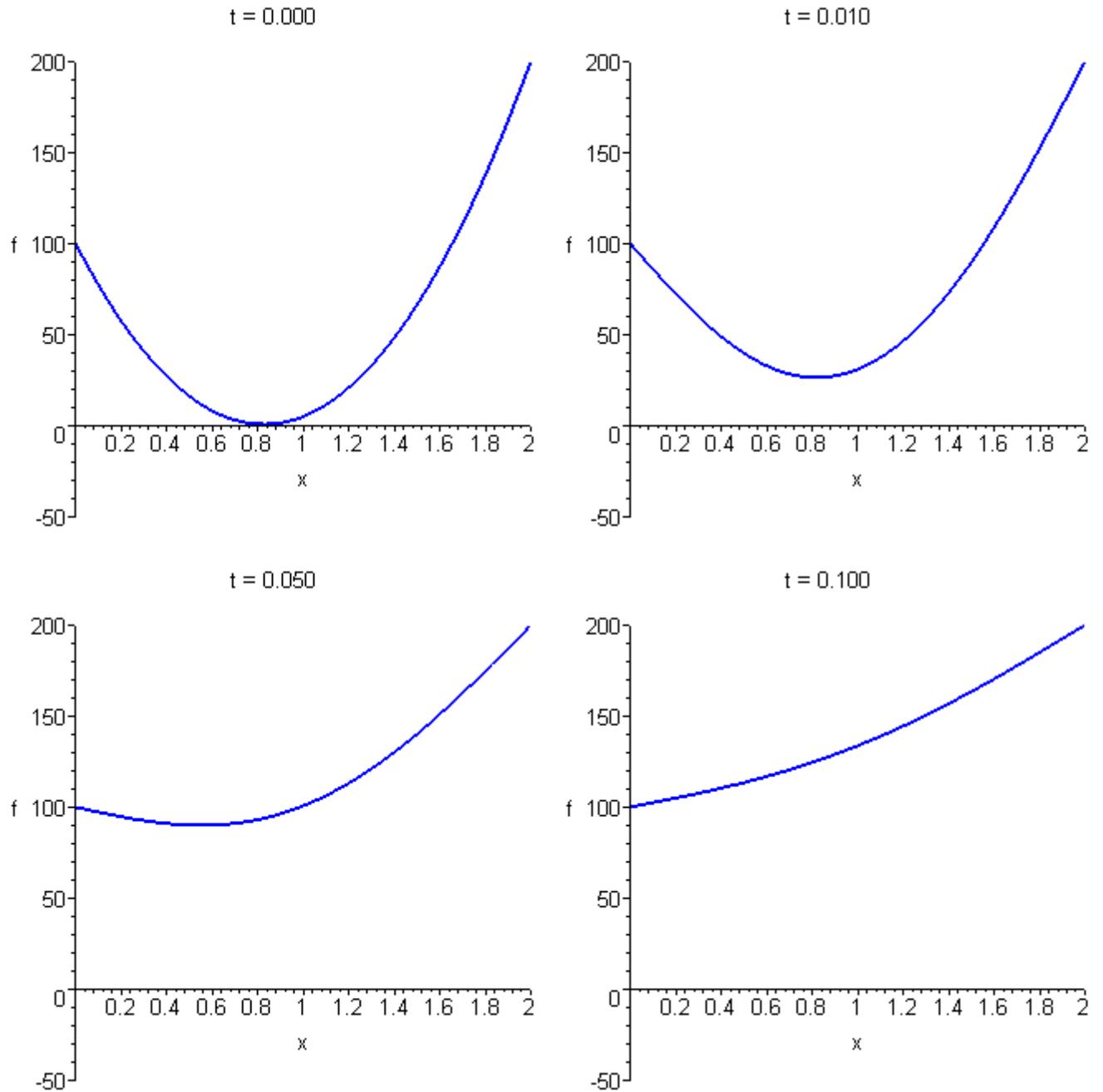
$$\Rightarrow c_n = 145 \left[\left((-z^2 + 2z) \frac{2}{n\pi} + \frac{16}{(n\pi)^3} \right) \cos\left(\frac{n\pi z}{2}\right) + \frac{8(z-1)}{(n\pi)^2} \sin\left(\frac{n\pi z}{2}\right) \right]_{z=0}^{z=2}$$

$$\Rightarrow c_n = \frac{2320}{(n\pi)^3} \left((-1)^n - 1 \right)$$

The complete solution is

$$u(x,t) = 50x + 100 - \frac{2320}{\pi^3} \sum_{n=1}^{\infty} \left(\frac{1 - (-1)^n}{n^3} \right) \sin\left(\frac{n\pi x}{2}\right) \exp\left(-\frac{9n^2\pi^2 t}{4}\right)$$

Some snapshots of the temperature distribution (from the tenth partial sum) from the Maple file at "www.engr.mun.ca/~ggeorge/9420/demos/ex8061.mws" are shown on the next page.

Example 8.06.1 (continued)

The steady state distribution is nearly attained in much less than a second!

END OF CHAPTER 8
END OF ENGI. 9420!
